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# Deformation mode-determined misorientation and microstructural characteristics in rolled pure Zr sheet

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In this work, commercially pure Zr sheets were subjected to  $\beta$  air cooling and then rolled to different reductions (10% and 50%) at room temperature. Microstructures of both the  $\beta$ -air-cooled and the rolled specimens were well characterized by electron channelling contrast imaging and electron backscatter diffraction techniques, with special attentions paid to their misorientation characteristics. Results show that the  $\beta$ -air-cooled specimen owns a Widmanstätten structure featured by lamellar grains with typical phase transformation misorientations. The 10% rolling allows prismatic slip and tensile twinning ( $\{11\cdot21\}\langle11\cdot2\cdot6\rangle$  and  $\{10\cdot12\}\langle10\cdot11\rangle$ ) to be activated profusely, which produce new low-angle ( $\sim 3^\circ$ – $5^\circ$ ) and high-angle ( $\sim 35^\circ$  and  $\sim 85^\circ$ ) misorientation peaks, respectively. After increasing the rolling reduction to 50%, twinning is suppressed and dislocation slip becomes the dominant deformation mode, with the lamellar grains highly elongated and aligned towards the rolling direction. Meanwhile, only one strong low-angle misorientation peak related to the prismatic slip is presented in the 50%-rolled specimen, with all other peaks disappeared. Analyses on local misorientations reveal that hardly any residual strains exist in the  $\beta$ -air-cooled specimen, which should be related to their sufficient relaxation during slow cooling. Residual strains introduced by 10% rolling are heterogeneously distributed near grain/twin boundaries while heavier deformation (50% rolling) produces much larger residual strains pervasively existing throughout the specimen microstructure.

**pure Zr, rolling, misorientation, twinning, electron backscatter diffraction**

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## 1 Introduction

Zr-based alloys have wide applications in nuclear, biomedical, and chemical industries because of attractive properties like low neutron absorption, excellent biocompatibility and good corrosion resistance [1–3]. During typical fabrications of most Zr alloy materials, a  $\beta$  cooling is often em-

ployed and followed by multiple-step hot/cold work and annealing [4–6]. To reach reliable property control, sufficient knowledge on microstructural characteristics evolved during such processings is necessarily required, with much already documented in literature by earlier researchers. Specifically, grain sizes and morphologies in many  $\beta$ -cooled Zr alloys with either martensite or Widmanstätten structures have been extensively studied [7–12]. Evolutions of structures, sizes and number densities of second phase particles in various newly developed Zr alloys are also the subject of plenty of

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investigations [13–18]. In recent years, characterization of deformed grain structures and their recrystallization behaviors of Zr materials have attracted a large number of attentions as well [19–24].

Misorientation has long been recognized to be another important microstructural parameter and can play a key role in influencing properties that depend on grain-boundary networks. In contrast with continuously active misorientation studies in cubic materials [25,26], misorientation characteristics in materials with a hexagonal close packed (hcp) structure have far less been explored [27,28]. The case is even worse for commercial Zr alloys which present hcp  $\alpha$  phase in most processing stages. Interestingly, recent work revealed that unique misorientation distributions would be presented in  $\beta$ -cooled Zr alloys [29,30]. How will such misorientation characteristic evolve during subsequent processings along with the underlying mechanism is also of great interest to be further explored. Such investigations are believed to be capable of not only deepening our understanding on misorientation features of hcp metals but also providing some direct insights into better microstructural control of Zr materials during fabrications.

Therefore, in the present work, commercially pure Zr sheets were selected as the experimental materials, which were  $\beta$  air-cooled and then rolled to different reductions (10% and 50%) at room temperature. Their microstructures were well investigated by use of electron channeling contrast (ECC) imaging and electron backscatter diffraction (EBSD) techniques based on a field emission gun scanning electron microscope, with particular attentions paid to the misorientation characteristics associated with specific deformation modes. Analyses based on local misorientations were also performed so as deformation-induced residual strains in various specimens could be evaluated.

## 2 Experimental

The as-received materials were 2-mm-thick commercially pure Zr (cp-Zr) sheets with major impurities of (in wt.%) of 1.15 Hf, <0.08 Fe+Cr and 0.15 O. The starting microstructure of the as-received sheets was noted to be featured by well-equiaxed grains after sufficient recrystallization, which has already been documented in our recent work [31]. The  $\beta$  transus temperature of the cp-Zr was determined to be about 960°C. The as-received materials were heat-treated at 1000°C for 10 min under vacuum in a box furnace and then cooled in flowing air. After the  $\beta$  air cooling, two sheets were rolled at room temperature by 5% per pass to thickness reductions of 10% and 50%, respectively. Rectangular specimens with dimensions of 12 and 10 mm along rolling and transverse directions (RD and TD) were cut from both the  $\beta$ -air-cooled and the rolled sheets.

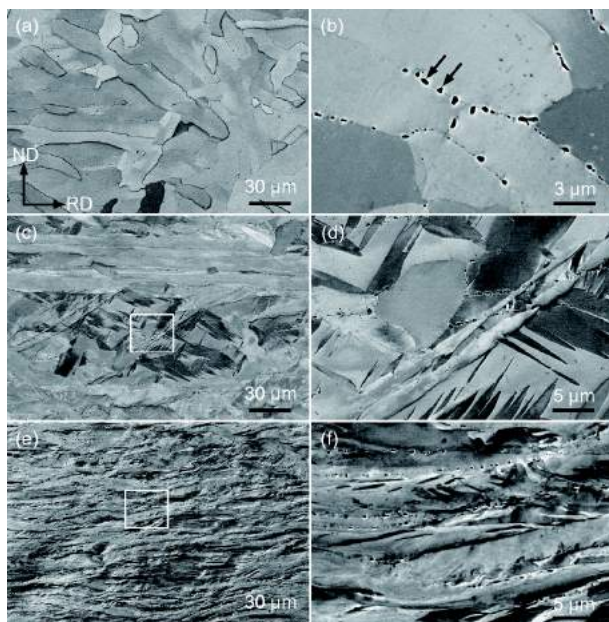
A Zeiss Sigma HD field emission gun scanning electron microscope (FEG-SEM) was employed to perform microstructural characterization for these specimens. ECC images based on back-scattered electrons were taken to reveal both composition- and orientation-related features [32–34]. An EBSD system attached to the FEG-SEM was used to further quantify crystallographic orientations of specific microstructures of the specimens. The EBSD system consisted of an Oxford Instruments NordlysMax2 detector with AZtec 2.4 and Channel 5 software packages for data acquisition and post processing, respectively. Prior to such microstructural characterization, RD-ND planes of the specimens were mechanically ground by SiC abrasive paper and then electropolished in a mixed solution of 10 mL perchloric acid, 20 mL butyl cellosolve and 70 mL methanol at 20 V and  $-30^{\circ}\text{C}$  for 30–60 s.

## 3 Results and discussion

Figure 1 presents direct microstructural observations (ECC images) for the  $\beta$ -air-cooled and the rolled specimens. After the  $\beta \rightarrow \alpha$  air cooling treatment, grains are found to be mainly presented in lamellar shape (Figure 1(a)), different from the equiaxed grains in the as-received material [31]. The average width of the lamellar grains is measured to be  $\sim 8.5\ \mu\text{m}$ . A magnified observation reveals that a number of second phase particles (SPPs) are distributed along boundaries of the lamellar grains, as arrowed in Figure 1(b). This suggests the occurrence of diffusion of impurity elements during the  $\beta$  air cooling and such microstructural features correspond to Widmanstätten structures. In fact, similar lamellar structures have been commonly seen in many  $\beta$ -treated Zr alloys and the SPPs at boundaries were determined to be Zr (Fe, Cr)<sub>2</sub> Laves phases [9,11]. Since the impurity content (especially Fe and Cr) in the cp-Zr is relatively limited, the volume fraction of the SPPs in the AC specimen is estimated to be very small (<0.5%).

After 10% rolling, most lamellar grains seem to be aligned and elongated slightly towards the RD, along with the presence of twin-like structures in some locations (Figure 1(c)). A closer observation in Figure 1(d) shows that the twin-like structures have distinct contrasts from the matrix with serrated boundaries between them. Sizes of the twin-like structures are found to vary a lot, ranging from several to near a hundred micrometers in length. For the 50%-rolled specimen, Figure 1(e) reveals that the prior Widmanstätten structures are heavily deformed and kinked, with the lamellar grains evidently elongated towards the RD. Grain refinement and subdivision could be noted after a magnified observation in Figure 1(f), as suggested by the greatly heterogeneous contrasts inside each lamellar grain.

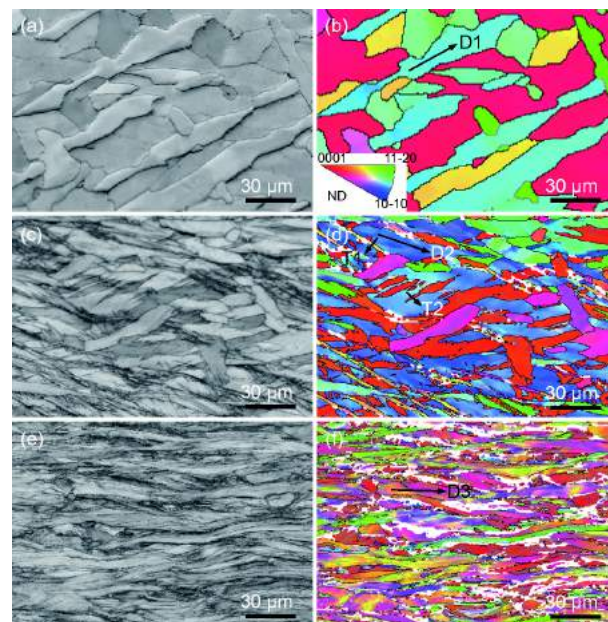
EBSD band contrast (BC) and inverse pole figure (IPF)



**Figure 1** ECC images of the ((a), (b))  $\beta$ -air-cooled, ((c), (d)) the 10%-rolled and ((e), (f)) the 50%-rolled specimens. Arrows in (b) indicate second phase particles; (d) and (f) correspond to boxed regions in (c) and (e), respectively.

maps of various specimens are displayed in **Figure 2**. **Figure 2(a)** reveals that high-quality Kikuchi patterns are obtained for the  $\beta$ -air-cooled specimen, suggesting hardly any residual stresses/strains. In **Figure 2(b)**, a rather uniform orientation (color) is noted for each lamellar grain enclosed by high angle boundaries (HABs,  $\theta > 15^\circ$ ). Very few low angle boundaries (LABs,  $2^\circ < \theta < 15^\circ$ ) can be found in the grain interior, confirming the suppression of dislocation activities during the  $\beta \rightarrow \alpha$  air cooling. Indeed, the structural discrepancy of  $\beta$  and  $\alpha$  phases might induce stresses during the allotropic phase transformation, leading to active slip and/or twinning like in the case of martensitic transformation [10,12]. In the present work, however, the relatively slow air cooling could allow such accumulated stresses to be sufficiently relaxed by creep processes at high temperature [35].

For the 10%-rolled specimen, varied image qualities in **Figure 2(c)** suggest the occurrence of heterogeneous plastic deformation. From **Figure 2(d)**, the twin-like structures revealed in **Figure 1(d)** can also be seen, with their orientations confirmed to be largely different from the matrix. In addition, in contrast with the insignificance of LABs in the  $\beta$ -air-cooled specimen, a moderate number of LABs are found to be dispersed inside both the matrix (lamellar grains) and the twin-like structures. As a result, the orientation (color) of a specific grain is no longer uniform. After increasing the rolling reduction to 50%, the heavily deformed structures revealed in **Figure 1(e)** are again noted in **Figure 2(e)**, from which the exclusive alignment towards the RD can be clearly seen for length directions of the lamellar grains. Moreover, **Figure 2(f)** reveals greatly denser LABs inside each grains



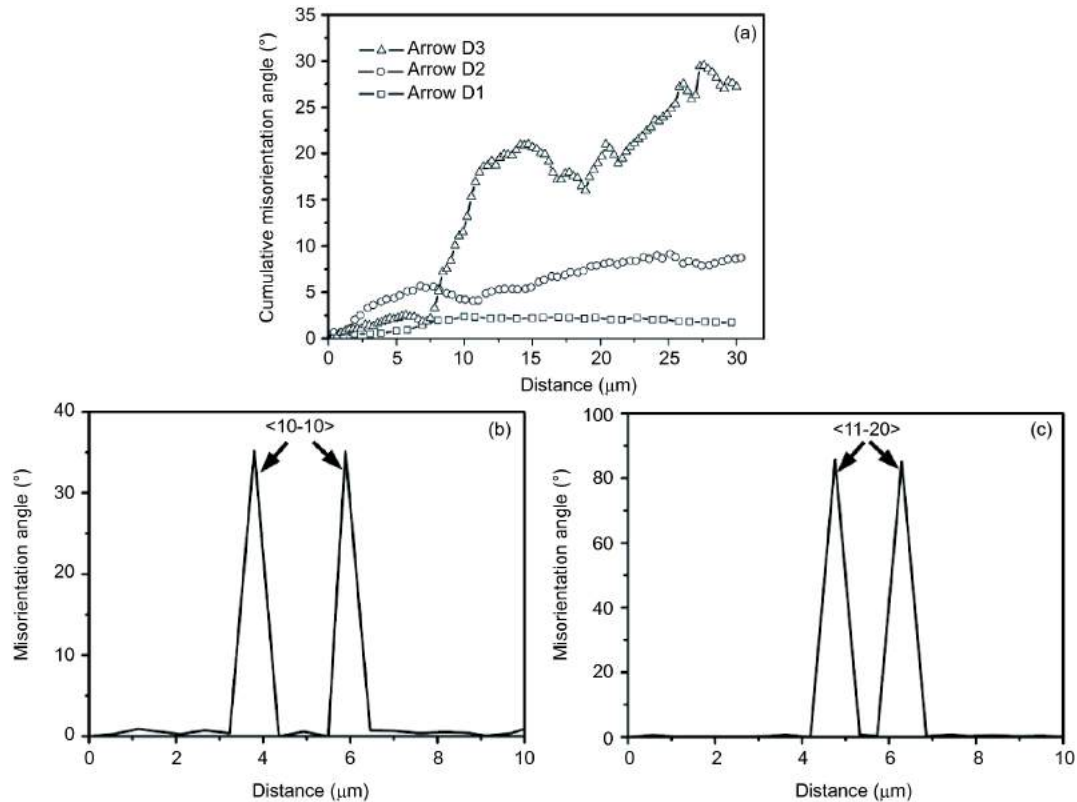
**Figure 2** (Color online) EBSD BC maps of (a) the  $\beta$ -air-cooled, (c) the 10%-rolled and (e) the 50%-rolled specimens; (b), (d) and (f) IPF maps corresponding accordingly to (a), (c) and (e), with HABs and LABs indicated by black and gray lines, respectively.

and much severer grain subdivision than in the 10%-rolled case. Also note that twin-like structures observed in the 10%-rolled specimen (**Figure 2(d)**) are no longer evident due probably to their morphological destruction by extensive dislocation movements.

Misorientation variations associated with typical microstructural characteristics in various specimens are presented in **Figure 3**. **Figure 3(a)** shows cumulative misorientation angles (CMAs) inside specific lamellar grains in different specimens. For the  $\beta$ -air-cooled specimen, the CMA is generally found to be low ( $< 2^\circ$ ), corresponding to a minimal orientation gradient ( $< 0.1^\circ \mu\text{m}^{-1}$ ). After a slight rolling (10% rolling), about  $9^\circ$  can be accumulated along a  $30 \mu\text{m}$  distance inside a lamellar grain, giving a moderate orientation gradient ( $\sim 0.3^\circ \mu\text{m}^{-1}$ ). Inside the lamellar grains of the 50%-rolled specimen, **Figure 3(a)** reveals that the CMA can be rapidly increased to near  $30^\circ$  for the same distance, indicating a large orientation gradient of  $\sim 1.0^\circ \mu\text{m}^{-1}$ . Note that developments of the orientation gradient during rolling are found to be consistent with increased densities of LABs (**Figure 2**). Also, the CMAs inside lamellar grains in both the rolled specimens are found to be gradually increased. One can thus deduce that the gradual lattice rotations during rolling should be induced by continuous dislocation slip rather than twinning [22,24].

With respect to twin-like structures appeared in the 10%-rolled specimen (**Figure 2(d)**), typical misorientation variations across their boundaries are also measured. As displayed in **Figure 3(b)** and (c), two types of misorientations (angle/axis pairs), i.e.,  $\sim 35^\circ / < 10-10 \rangle$  and  $\sim 85^\circ / < 11-20 \rangle$ , are gen-





**Figure 3** Cumulative misorientation angles along various measuring paths as arrowed in Figure 2. (a) D1-D3; (b) T1; (c) T2.

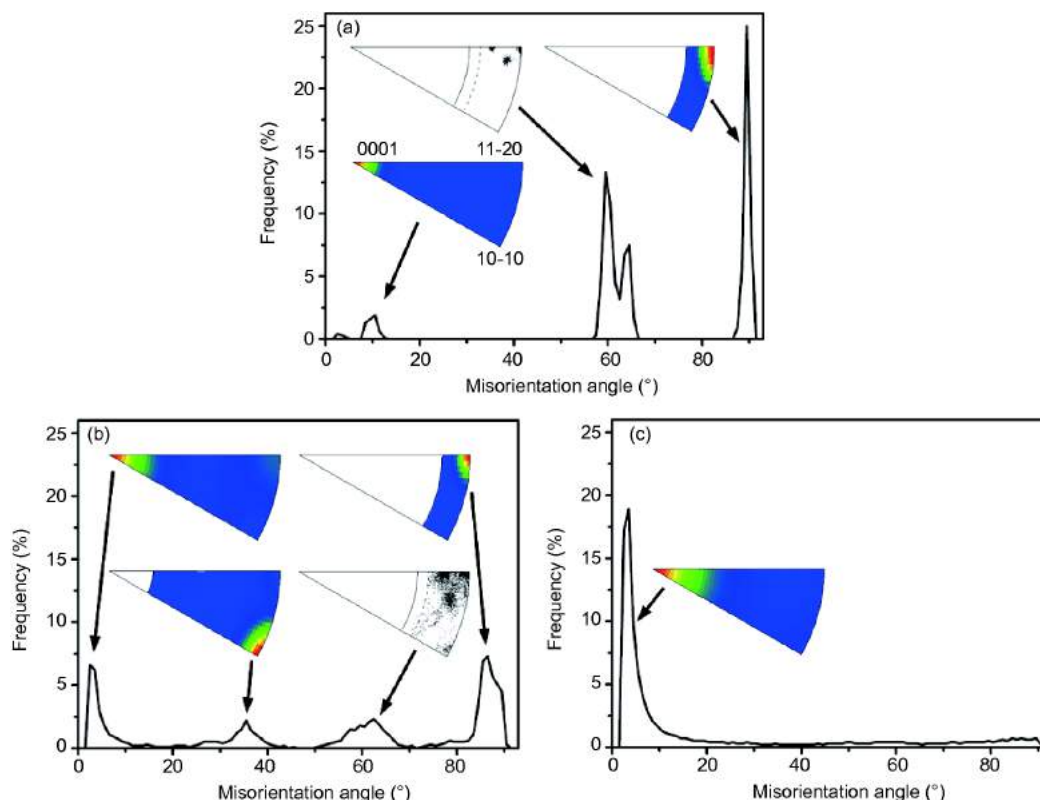
erally detected. They are found to be well coincident with crystallographic characteristics of  $\{11\text{-}21\}\langle 11\text{-}2\text{-}6\rangle$  and  $\{10\text{-}12\}\langle 10\text{-}11\rangle$  tensile twins (Table 1). It is thus known that such twin-like structures correspond to the two types of tensile twins in  $\alpha$ -Zr. In an earlier work, Crépin et al. [36] also detected the two types of tensile twins in a  $\beta$ -treated Zr702 subjected to uniaxial tension. They reported that the measured critical resolved shear stress (CRSS) of the  $\{10\text{-}12\}\langle 10\text{-}11\rangle$  twins is much lower than that of the  $\{11\text{-}21\}\langle 11\text{-}2\text{-}6\rangle$  twins and thus allowed them to be more easily initiated.

In addition to the local misorientation characteristics revealed in Figure 3, Figure 4 presents global misorientation

angle and rotation axis distributions of all the three specimens. For the  $\beta$ -air-cooled specimen, three peaks around  $10^\circ$ ,  $60^\circ$  and  $90^\circ$  can be exclusively seen from Figure 4(a). Such a misorientation angle distribution has been commonly observed for Zr alloys after  $\beta$ -cooling [10,30,37]. In fact, the classical Burgers orientation relationship (OR) ( $\{110\}_\beta//\{0001\}_\alpha$  and  $\langle 111\rangle_\beta//\langle 11\text{-}20\rangle_\alpha$ ) predicts five specific misorientations between  $\alpha$  variants generated by the same  $\beta$  orientation, i.e.,  $10.5^\circ/\langle 0001\rangle$ ,  $60^\circ/\langle 11\text{-}20\rangle$ ,  $60.8^\circ/\langle 12\text{-}31\rangle$ ,  $63.3^\circ/\langle 44\text{-}83\rangle$  and  $90^\circ/\langle 12\text{-}30\rangle$  [30]. Examinations on rotation axes corresponding to the three peaks suggest that they are in good agreement with the five misorientations predicted by the Burgers OR. The peak around  $60^\circ$  can be

**Table 1** Deformation modes operable in  $\alpha$ -Zr and their basic parameters.

| Deformation mode |             |   | Taylor rotation axis                   | Lattice rotation | Twinning shear |
|------------------|-------------|---|--|------------------|----------------|
| Twinning         | Tensile     | $\{10\text{-}12\}\langle 10\text{-}11\rangle$         | $\langle 11\text{-}20\rangle$          | $85.22^\circ$    | 0.167          |
|                  |             | $\{11\text{-}21\}\langle 11\text{-}2\text{-}6\rangle$ | $\langle 10\text{-}10\rangle$          | $34.84^\circ$    | 0.627          |
|                  | Compressive | $\{11\text{-}22\}\langle 11\text{-}2\text{-}3\rangle$ | $\langle 10\text{-}10\rangle$          | $64.22^\circ$    | 0.225          |
|                  |             | $\{10\text{-}11\}\langle 10\text{-}12\rangle$         | $\langle 11\text{-}20\rangle$          | $57.05^\circ$    | 0.104          |
| Slip             | <a> type    | $\{10\text{-}10\}\langle 11\text{-}20\rangle$         | $\langle 0001\rangle$                  | Continuous       | Not applicable |
|                  |             | $\{0002\}\langle 11\text{-}20\rangle$                 | $\langle 10\text{-}10\rangle$          |                  |                |
|                  |             | $\{10\text{-}11\}\langle 11\text{-}20\rangle$         | $\langle 10\text{-}12\rangle$          |                  |                |
|                  | <c+a> type  | $\{10\text{-}11\}\langle 11\text{-}23\rangle$         | $\langle 13\text{-}8\text{-}53\rangle$ |                  |                |
|                  |             | $\{11\text{-}21\}\langle 11\text{-}23\rangle$         | $\langle 6\text{-}1\text{-}53\rangle$  |                  |                |



**Figure 4** (Color online) Misorientation angle and rotation axis distributions of (a) the  $\beta$ -air-cooled, (b) the 10%-rolled and (c) the 50%-rolled specimens.

ascribed to an overlap of measured angles corresponding to the second, third and fourth Burgers misorientations (see ref. [30] for detailed illustration). One can know that the misorientation characteristic of the  $\beta$ -air-cooled specimen is completely resulted from the  $\beta \rightarrow \alpha$  phase transformation according to the Burgers OR.

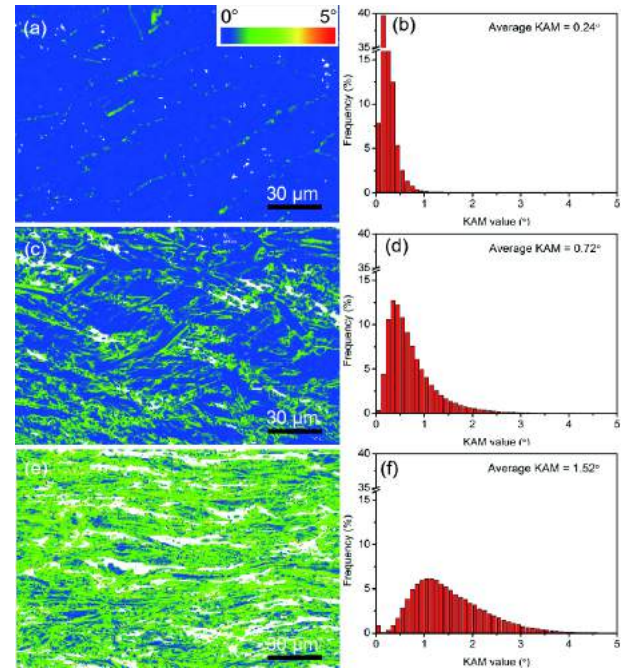
For the 10%-rolled specimen, remarkable changes can be noted for its misorientation angle distribution (Figure 4(b)). The highest peak near  $90^\circ$  in the  $\beta$ -air-cooled specimen is largely reduced in intensity and slightly shifted towards lower angles. In fact, its new peak center is adjusted close to  $85^\circ$  with the rotation axis well concentrated on  $\langle 11-20 \rangle$ . This suggests that the prior  $90^\circ/\sim \langle 12-30 \rangle$  misorientation should have been largely destructed and overlapped with the  $\sim 85^\circ/\langle 11-20 \rangle$  misorientation due to the appearance of dense  $\{10-12\}\langle 10-11 \rangle$  twins (Figure 2(d)). The prior peak around  $60^\circ$  is also heavily decayed, with evidently spread angles and rotation axes. A new peak around  $35^\circ$  appears with rotation axes focused on  $\langle 10-10 \rangle$ , corresponding to that of  $\{11-21\}\langle 11-2-6 \rangle$  twins. The presence of this peak also suggests the abundance of the  $\{11-21\}\langle 11-2-6 \rangle$  twins in the 10%-rolled specimen. This is in contrast with the case that no twins were found in a typically textured pure Zr sheet subjected to the same rolling [22], in spite of its larger grain sizes ( $\sim 9.6 \mu\text{m}$  on average). This may be mainly attributed to a textural difference between them. Indeed, the  $\beta$ -air-cooled specimen

in this work was found to possess a relatively random texture (documented in ref. [7]), which would contain more favorable orientations for twinning. With respect to low misorientation angles, the prior peak near  $\sim 10^\circ$  seems to be completely suppressed and replaced by a new one focused on lower angles ( $3^\circ\text{--}5^\circ$ ). The new low-angle peak should mainly correspond to those LABs produced by dislocation movements, as shown in Figure 2(d). According to Table 1, activated slip modes leading to these low angles can be probed by considering their corresponding rotation axes (more detailed illustration can be seen in ref. [38]). As revealed in Figure 4(b), a highly concentrated rotation axis of  $\langle 0001 \rangle$  is found for these low angles, suggesting prismatic slip ( $\{10-10\}\langle 11-20 \rangle$ ) of  $\langle a \rangle$  dislocation as the main slip mode. In fact, compared with other operable slip systems, prismatic slip is usually the easiest to be initiated during deformation due to its much lower CRSS [39].

With increasing rolling reduction to 50%, all the prior high-angle peaks appear to be completely suppressed, with only one low-angle peak highly intensified (Figure 4(c)). Clearly, twinning should no longer play a key role in the deformation process, which could reply more on slip instead. A similar phenomenon was earlier reported for a Zircaloy-2 subjected to different degrees of uniaxial compression [40]. The insignificance of twinning at later stages of deformation can be ascribed to both changed crystallographic texture and

refined grains by earlier twins. Figure 4(c) also reveals that the rotation axes corresponding to the left low-angle peak remain concentrated on the  $\langle 0001 \rangle$ , indicating the same predominant deformation mode (prismatic slip) as that in 10%-rolled specimen. Nevertheless, since the prismatic slip could only accommodate strains along a-axes of hcp structures, other deformation modes with the ability of accommodating strains along grain c-axes should be needed. Such a role could be played by twinning at the early state of rolling (Figures 2(d) and 4(b)). However, the absence of twins in the 50%-rolled specimen suggests that pyramidal slip of  $\langle c+a \rangle$  dislocations should have been activated to some extent, in spite of their higher CRSSs [39]. This is also supported by the earlier work [41] that presented experimental evidences of activation of  $\{10\text{-}11\}\langle 11\text{-}23 \rangle$  slip in a pure Zr deformed at room temperature.

Figure 5 presents EBSD kernel average misorientation (KAM) maps, which are reconstructed in terms of local misorientation characteristics and known to be capable of directly reflect distributions of residual strains in the specimen interiors [42]. For each KAM map, the maximum misorientation is selected to be  $5^\circ$  with only the three nearest neighbors considered. Figure 5(a) reveals a relatively low and uniform KAM distribution throughout the microstructure of the  $\beta$ -air-cooled specimen, with an average KAM calculated to be  $0.24^\circ$  (Figure 5(b)). The value is comparable to that of fully recrystallized grains [43], suggesting few residual strains after the slow  $\beta \rightarrow \alpha$  cooling, consistent with the above analysis based on Kikuchi patterns (Figure 2(a)). As above mentioned, this could be ascribed to sufficient relaxation of stresses/strains by creep at high temperature [35]. For the 10%-rolled specimen, the KAM distribution is found to become highly heterogeneous in Figure 5(c). A comparison between Figures 2(d) and 5(c) suggests that the KAM is generally larger near grain/twin boundaries than inside them. This can be understood that dislocations are easily to be piled up in front of the grain or twin boundaries (both of them belong to HABs) during the rolling. The increased average KAM ( $0.72^\circ$ ) and the more spread distribution shown in Figure 5(d) clearly manifest the presence of residual strains. After the heavier deformation (50% rolling), larger KAM can be noticed throughout the specimen microstructure, as shown in Figure 5(e). The average KAM of the 50%-rolled specimen is calculated to be  $1.52^\circ$  (more than double that of the 10%-rolled specimen) with its distribution histogram spreading more towards higher angles (Figure 5(f)). According to the KAM results, much larger residual strains can be expected for the 50%-rolled specimen. As twinning is well suppressed, the residual strains should be essentially attributed to extensive dislocation slip (both prismatic and pyramidal systems). Also, different from the preferred activities of dislocations near HABs during the 10% rolling (Figure 5(c)), the pervasive



**Figure 5** (Color online) KAM maps of (a) the  $\beta$ -air-cooled, (c) the 10%-rolled and (e) the 50%-rolled specimens; (b), (d) and (f) KAM distribution histograms corresponding accordingly to (a), (c) and (e).

existence of residual strains (Figure 5(e)) indicates active dislocation slip both inside grains and near their boundaries. This can also be evidenced by the large orientation gradient in lamellar grain interiors, as revealed in Figure 3(a).

## 4 Conclusions

(1) The  $\beta$ -air-cooled specimen possesses a Widmanstätten structure featured by lamellar grains with scattered orientations. During the 10% rolling, both active twinning and slip occur; however, only dislocation slip is the dominant deformation mode after increasing the rolling reduction to 50%, with the lamellar grains highly elongated and aligned towards the RD.

(2) Three typical misorientation angle peaks corresponding to the Burgers misorientations are noted for the  $\beta$ -air-cooled specimen. After the 10% rolling, all the three peaks are largely decayed with the appearance of new low-angle ( $\sim 3^\circ$ – $5^\circ$ ) and high-angle ( $\sim 35^\circ$  and  $\sim 85^\circ$ ) peaks, resulting from prismatic slip and tensile twinning ( $\{11\text{-}21\}\langle 11\text{-}2\text{-}6 \rangle$  and  $\{10\text{-}12\}\langle 10\text{-}11 \rangle$ ), respectively. Only one strong low-angle peak related to prismatic slip is presented after increasing the rolling reduction to 50%, with all other peaks disappeared.

(3) Hardly any residual strains are found in the  $\beta$ -air-cooled specimen, which should be related to sufficient relaxation during slow cooling. Residual strains introduced by 10% rolling are heterogeneously distributed near grain/twin

boundaries due to preferred dislocation pile ups. After heavier deformation (50% rolling), much larger residual strains are noted with the pervasive existence throughout the specimen microstructure.

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